

Consider the infinitesimal line element of  $n$  space, Developed Version

$$ds^2 = c^2 d\tau^2 = c^2 m(r) dt^2 - \frac{dr^2}{m(r)} - r^2 d\phi^2 \quad (1)$$

and define the linear velocity as:

$$v^2 dt^2 = \frac{dr^2}{m(r)} + r^2 d\phi^2 \quad (2)$$

Since  $m(r)$  is considered to be independent of time:

$$\begin{aligned} v^2 &= \frac{1}{m(r)} \left( \frac{dr}{dt} \right)^2 + r^2 \left( \frac{d\phi}{dt} \right)^2 \\ &= \frac{1}{m(r)} \dot{r}^2 + r^2 \dot{\phi}^2 \end{aligned} \quad (3)$$

It follows that:

$$c^2 d\tau^2 = c^2 m(r) dt^2 - v^2 dt^2 \quad (4)$$

So

$$\gamma = \frac{dt}{d\tau} = \left( m(r) - \frac{v^2}{c^2} \right)^{-1/2} \quad (5)$$

is the Lorentz factor in  $n$  space.

The structure of the field equations of ECE2 theory is the structure of the field equations in Minkowski spacetime. However ECE2 is generally covariant and developed in a space with finite torsion and curvature. It is therefore considered that the linear velocity in  $n$  theory is:

$$\underline{v} = \gamma \underline{\dot{r}} \quad (6)$$

From eq. (3):

$$v^2 = \underline{\dot{r}} \cdot \underline{\dot{r}} = \frac{1}{m(r)} \dot{r}^2 + r^2 \dot{\phi}^2 \quad (7)$$

So

$$\underline{\dot{r}} = \frac{\dot{r}}{m(r)^{1/2}} \underline{e}_r + r\dot{\phi} \underline{e}_\phi \quad - (8)$$

where  $\underline{e}_r$  and  $\underline{e}_\phi$  are the unit vectors of the plane polar system.

It follows that:

$$\begin{aligned} \underline{\ddot{r}} &= \frac{1}{m(r)^{1/2}} \frac{d}{dt} (\dot{r} \underline{e}_r) + \frac{d}{dt} (r\dot{\phi} \underline{e}_\phi) \\ &= \frac{1}{m(r)^{1/2}} (\ddot{r} \underline{e}_r + \dot{r} \dot{\underline{e}}_r) + \dot{r}\dot{\phi} \underline{e}_\phi + r\ddot{\phi} \underline{e}_\phi + r\dot{\phi} \dot{\underline{e}}_\phi \\ &= \frac{1}{m(r)^{1/2}} (\ddot{r} \underline{e}_r + \dot{r}\dot{\phi} \underline{e}_\phi) + (\dot{r}\dot{\phi} + r\ddot{\phi}) \underline{e}_\phi - r\dot{\phi}^2 \underline{e}_r \\ &= \left( \frac{\ddot{r}}{m(r)^{1/2}} - r\dot{\phi}^2 \right) \underline{e}_r + \left( r\ddot{\phi} + \dot{r}\dot{\phi} \left( 1 + \frac{1}{m(r)^{1/2}} \right) \right) \underline{e}_\phi \end{aligned}$$

The acceleration due to gravity is: - (9)

$$\begin{aligned} \underline{g} &= \frac{d}{dt} (\gamma \underline{\dot{r}}) \\ &= \frac{d\gamma}{dt} \underline{\dot{r}} + \gamma \underline{\ddot{r}} \end{aligned} \quad - (10)$$

so in GR theory:

$$\underline{g} = \frac{d\gamma}{dt} \underline{\dot{r}} + \gamma \underline{\ddot{r}} = -\underline{\nabla} \underline{\Phi} + \underline{\Omega} \underline{\Phi} \quad - (11)$$

$$\underline{\Phi} = -\frac{MG}{r} \quad - (12)$$

of gravitational potential.  
It follows that:

$$\begin{aligned} \underline{g} &= \frac{dY}{dt} \left( \frac{\dot{r}}{m(r)^{1/2}} \underline{e}_r + r\dot{\phi} \underline{e}_\phi \right) \\ &+ \gamma \left( \left( \frac{\ddot{r}}{m(r)^{1/2}} - r\dot{\phi}^2 \right) \underline{e}_r + \left( r\ddot{\phi} + \dot{r}\dot{\phi} \left( 1 + \frac{1}{m(r)^{1/2}} \right) \right) \underline{e}_\phi \right) \\ &= \left( -\frac{mG}{r^2} + \Omega_r \underline{\Phi} \right) \underline{e}_r \quad - (13) \\ &= -\frac{mG}{r} \left( \frac{1}{r} + \Omega_r \right) \underline{e}_r \end{aligned}$$

Here it has been assumed that:

$$\underline{\Omega} = \Omega_r \underline{e}_r + \Omega_\phi \underline{e}_\phi \quad - (14)$$

and that

$$\Omega_\phi = 0 \quad - (15)$$

From eq. (13) - (15):

$$\frac{dY}{dt} \frac{\dot{r}}{m(r)^{1/2}} + \gamma \left( \frac{\ddot{r}}{m(r)^{1/2}} - r\dot{\phi}^2 \right) = -\frac{mG}{r} \left( \frac{1}{r} + \Omega_r \right) \quad - (16)$$

is the generally covariant Leibniz equation in n space.

$$\text{and} \quad \frac{dL}{dt} = r\dot{\phi} \frac{dY}{dt} + \gamma \left( r\ddot{\phi} + \dot{r}\dot{\phi} \left( 1 + \frac{1}{m(r)^{1/2}} \right) \right) = 0 \quad - (17)$$

is the equation of conservation of angular momentum L.  
 The result is given by solving eqs. (16) and (17) simultaneously, and L is found by integrating eq. (17) numerically.

i) As in UFT 190, the orbit can also be found by numerical integration of:

$$\frac{dr}{d\phi} = r^2 \left( \frac{1}{b^2} - m(r) \left( \frac{1}{a^2} + \frac{1}{r^2} \right) \right)^{1/2} \quad (18)$$

and in ETC cosmology:

$$m(r) = 2 - \exp \left( 2 \exp \left( -\frac{r}{R} \right) \right) \quad (19)$$

where  $R$  is a characteristic distance in the universe.

In papers such as UFT 108 and UFT 190, it was shown that orbital shrinkage can be explained with n theory.

The Lagrangian of the above theory is found from:

$$\underline{p} = m \underline{v} = m \gamma \underline{\dot{r}} = \frac{\partial \mathcal{L}}{\partial \underline{\dot{r}}} \quad (20)$$

where:

$$\gamma = \frac{dt}{d\tau} = \left( m(r) - \frac{v^2}{c^2} \right)^{-1/2} \quad (21)$$

$$= \left( m(r) - \frac{\dot{\underline{r}} \cdot \dot{\underline{r}}}{c^2} \right)^{-1/2}$$

This guarantees that the Lagrangian and metric theories give the same orbit equations (16) and (17).