

# 10(4): Origin and Derivation of the Precession due to FIEA Covariant General Relativity

Consider the invariance of phase under four rotation:

$$\phi = \kappa^\mu x_\mu = \kappa'^\mu x'_\mu \quad (1)$$

$$= \omega t - \underline{\kappa} \cdot \underline{r} = \omega' t' - \underline{\kappa}' \cdot \underline{r}'$$

Here:  $\kappa^\mu = \left( \frac{\omega}{c}, \underline{\kappa} \right), x^\mu = (ct, \underline{r}) \quad (2)$

In its infinitesimal format:

$$\phi = \omega dt - \underline{\kappa} \cdot d\underline{r} = \omega' dt' - \underline{\kappa}' \cdot d\underline{r}' \quad (3)$$

A particle is moving frame  $\kappa'$  does not move w.r.t respect to  $\kappa'$  so:  $d\underline{r}' = 0 \quad (4)$

and  $\phi = \omega dt - \underline{\kappa} \cdot d\underline{r} = \omega' dt' \quad (5)$

In the notation of previous notes:

$$\phi = \omega dt - \underline{\kappa} \cdot d\underline{r} = \omega_0 dt_1 = \omega_0 d\tau \quad (6)$$

where  $\omega_0$  is the rest frequency and  $\tau$  the proper time.

Now use the de Broglie / Einstein equations:

$$E_h = \hbar \omega = \gamma mc^2 \quad (7)$$

$$p = \hbar \underline{\kappa} = \gamma m \underline{v} \quad (8)$$

to find that the rest angular frequency is:

$$\omega_0 = \frac{mc^2}{\hbar} \quad (9)$$

2) This is the angular frequency in frame  $K'$ . In frame  $K$ , the laboratory or observer frame:

$$\omega = \gamma \omega_0 \quad (10)$$

In the preceding notes it has been shown that:

$$dt = \gamma dt_1 \quad (11)$$

It follows that the phase change between frame  $K$  and frame  $K'$  is:

$$\begin{aligned} \Delta \phi &= \omega t - \omega_0 t_1 = \underline{\kappa} \cdot \underline{r} \\ &= \omega t - \frac{\omega t}{\gamma} = \omega t \left(1 - \frac{1}{\gamma}\right) \quad (12) \end{aligned}$$

For one revolution:  $\omega t = 2\pi \quad (13)$

$$\text{so } \Delta \phi = 2\pi \left(1 - \frac{1}{\gamma}\right) = \underline{\kappa} \cdot \underline{r} \quad (14)$$

$$= 2\pi \frac{(\gamma^2 - 1)}{\gamma^2} = 2\pi (\gamma - 1) \frac{(\gamma + 1)}{\gamma^2}$$

The definition of precession used in the "usual" theory of Thomas precession is:

$$\Delta \phi = 2\pi (\gamma - 1) = \frac{\gamma^2}{\gamma + 1} \underline{\kappa} \cdot \underline{r} \quad (15)$$

The correct expression is:

$$\Delta \phi = \omega t - \omega_0 t_1 = 2\pi \left(1 - \frac{1}{\gamma}\right) = \underline{\kappa} \cdot \underline{r} \quad (16)$$

In the low velocity limit eq. (16) reduces to:

$$\Delta\phi = 2\pi \left( 1 - \frac{1}{\gamma} \right) = 2\pi \frac{v^2}{c^2} \quad - (17)$$

and this is true for all  $\gamma$  and all  $v$ . Eq. (17) is the best result of fundamental theory, and replace the Sittery "Wikipedia" result:

$$\Delta\phi = ? 2\pi(\gamma - 1) \quad - (18)$$

Eq. (17) is an exact result. If de Sitter present:

$$v^2 = v_N^2 + 3\omega^2 r^2 \quad - (19)$$

and the universal law of precession is:

$$\Delta\phi = \omega t - \omega_0 t_1 = 2\pi \frac{v^2}{c^2} = 2\pi \frac{(v_N^2 + 3\omega^2 r^2)}{c^2} \quad - (20)$$

If de Sitter rotation is absent:

$$\Delta\phi = \omega t - \omega_0 \tau = 2\pi \frac{v_N^2}{c^2} \quad - (21)$$

and this is precession due to time dilatation and length contraction, in which

$$\gamma = \frac{l'}{l} = \frac{\Delta t}{\Delta\tau} \quad - (22)$$

$$\Delta\phi = 2\pi \left( 1 - \frac{l^2}{l'^2} \right) = 2\pi \left( 1 - \left( \frac{\Delta\tau}{\Delta t} \right)^2 \right) \quad - (23)$$