

23(8): Final Expression for the Electric Field strength \underline{E} .

The dipole electric field strength in the presence of the vacuum is:

$$\underline{E} = \frac{1}{4\pi\epsilon_0 r^5} (\underline{r} + \delta\underline{r}) (\underline{p} \cdot (\underline{r} + \delta\underline{r})) \left(1 - \frac{5x}{2}\right) - \frac{p}{4\pi\epsilon_0 r^3} \left(1 - \frac{3x}{2}\right) \quad (1)$$

also $x = \frac{1}{r^2} (2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r}) \quad (2)$

So:

$$\underline{E} = \frac{1}{4\pi\epsilon_0 r^5} \left((\underline{r} + \delta\underline{r}) \underline{p} \cdot (\underline{r} + \delta\underline{r}) - \frac{5x}{2} (\underline{r} + \delta\underline{r}) (\underline{p} \cdot (\underline{r} + \delta\underline{r})) \right) - \frac{p}{4\pi\epsilon_0 r^3} + \frac{3x p}{8\pi\epsilon_0 r^3}$$

$$= \frac{1}{4\pi\epsilon_0 r^5} (\underline{r} \underline{p} \cdot \underline{r} + \delta\underline{r} \underline{p} \cdot \underline{r} + \underline{r} \underline{p} \cdot \delta\underline{r} + \delta\underline{r} \underline{p} \cdot \delta\underline{r}) \left(1 - \frac{5x}{2}\right) - \frac{p}{4\pi\epsilon_0 r^3} \left(1 - \frac{3x}{2}\right) \quad (2)$$

$$= \frac{1}{4\pi\epsilon_0 r^3} \left(\frac{\underline{r} \underline{p} \cdot \underline{r}}{r^2} - \underline{p} \right) - \frac{x}{8\pi\epsilon_0 r^3} \left(\frac{5 \underline{r} \underline{p} \cdot \underline{r}}{r^2} - 3 \underline{p} \right) + \frac{1}{4\pi\epsilon_0 r^5} (\delta\underline{r} \underline{p} \cdot \underline{r} + \underline{r} \underline{p} \cdot \delta\underline{r} + \delta\underline{r} \underline{p} \cdot \delta\underline{r}) \left(1 - \frac{5x}{2}\right)$$

As shown in Note 393(7):

$$\begin{aligned} \langle x \rangle &= \quad (3) \\ &= \frac{1}{r^2} \langle 2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r} \rangle \\ &= \frac{1}{r^2} \langle \delta\underline{r} \cdot \delta\underline{r} \rangle \end{aligned}$$

in an isotropic sample, where:

$$\langle \underline{\delta r} \rangle = \underline{0} \quad - (4)$$

Therefore:

$$\langle \underline{E} \rangle = \underline{E}_0 - \frac{\langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{8\pi\epsilon_0 r^5} \left(\frac{5 \underline{r} \underline{p} \cdot \underline{r}}{r^2} - 3 \underline{p} \right) + \frac{1}{4\pi\epsilon_0 r^5} \left\langle \underline{\delta r} \underline{p} \cdot \underline{r} + \underline{r} \underline{p} \cdot \underline{\delta r} + \underline{\delta r} \underline{p} \cdot \underline{\delta r} \left(1 - \frac{5x}{2} \right) \right\rangle \quad - (5)$$

where

$$\underline{E}_0 = \frac{1}{4\pi\epsilon_0 r^3} \left(\frac{\underline{r} \underline{p} \cdot \underline{r}}{r^2} - \underline{p} \right) \quad - (6)$$

is the standard model dipole field, the field in the hypothetical absence of the vacuum.

From observations of the Lamb shift it is known that:

$$\frac{5x}{2} \ll (1 - \dots) \quad - (8)$$

So:

$$\langle \underline{E}_1 \rangle = \frac{1}{4\pi\epsilon_0 r^5} \left(\langle \underline{\delta r} \underline{p} \cdot \underline{r} \rangle + \langle \underline{r} \underline{p} \cdot \underline{\delta r} \rangle + \langle \underline{\delta r} \underline{p} \cdot \underline{\delta r} \rangle \right)$$

to an excellent approximation

By isotropy:

$$\langle \underline{\delta r} \underline{p} \cdot \underline{r} \rangle = \langle \underline{r} \underline{p} \cdot \underline{\delta r} \rangle = \underline{0} \quad - (9)$$

3) so:

$$\langle \underline{E}_1 \rangle = \frac{1}{4\pi\epsilon_0 r^5} \langle \underline{\delta r} \underline{p} \cdot \underline{\delta r} \rangle - (10)$$

$$= \frac{1}{4\pi\epsilon_0 r^5} \langle (p_x \delta x + p_y \delta y + p_z \delta z) (\delta x \underline{i} + \delta y \underline{j} + \delta z \underline{k}) \rangle$$

By isotropy:

$$\langle \delta x \delta y \rangle = \langle \delta x \delta z \rangle = \langle \delta y \delta z \rangle = 0 - (11)$$

so

$$\begin{aligned} \langle \underline{E}_1 \rangle &= \frac{1}{4\pi\epsilon_0 r^5} \langle p_x \delta x^2 \underline{i} + p_y \delta y^2 \underline{j} + p_z \delta z^2 \underline{k} \rangle \\ &= \frac{1}{4\pi\epsilon_0 r^5} \left(p_x \underline{i} \langle \delta x^2 \rangle + p_y \underline{j} \langle \delta y^2 \rangle + p_z \underline{k} \langle \delta z^2 \rangle \right) - (12) \end{aligned}$$

By isotropy:

$$\langle \delta x^2 \rangle = \langle \delta y^2 \rangle = \langle \delta z^2 \rangle - (13)$$

and

$$\langle \delta r^2 \rangle = \langle \delta x^2 \rangle + \langle \delta y^2 \rangle + \langle \delta z^2 \rangle - (14)$$

so

$$\langle \delta r^2 \rangle = 3 \langle \delta x^2 \rangle = 3 \langle \delta y^2 \rangle = 3 \langle \delta z^2 \rangle - (15)$$

and

$$\begin{aligned} \langle \delta x^2 \rangle &= \langle \delta y^2 \rangle = \langle \delta z^2 \rangle = \frac{1}{3} \langle \delta r^2 \rangle \\ &= \frac{1}{3} \langle \underline{\delta r} \cdot \underline{\delta r} \rangle - (16) \end{aligned}$$

It follows that:

$$\langle \underline{E}_1 \rangle = \frac{\langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{4\pi\epsilon_0 r^5} \underline{p} - (17)$$

4) Finally:

$$\langle \underline{E} \rangle = \underline{E}_0 - \frac{\langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{8\pi\epsilon_0 r^5} \left(\frac{5\underline{r} \underline{p} \cdot \underline{r}}{r^2} - 3\underline{p} \right) + \frac{\langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{4\pi\epsilon_0 r^5} \underline{p} \quad - (18)$$

$$\langle \underline{E} \rangle = \underline{E}_0 - \frac{5 \langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{8\pi\epsilon_0 r^5} \left(\frac{\underline{r} \underline{p} \cdot \underline{r}}{r^2} - \underline{p} \right) \quad - (19)$$

It is seen that the correction due to the vacuum is proportional to $\langle \underline{\delta r} \cdot \underline{\delta r} \rangle$.

If the approximation (7) is not used, other terms appear in eq. (19), and these are worked out in Section 3 of UFT593.