

69(1): Motion of a Gyroscope in a Gravitational Field.

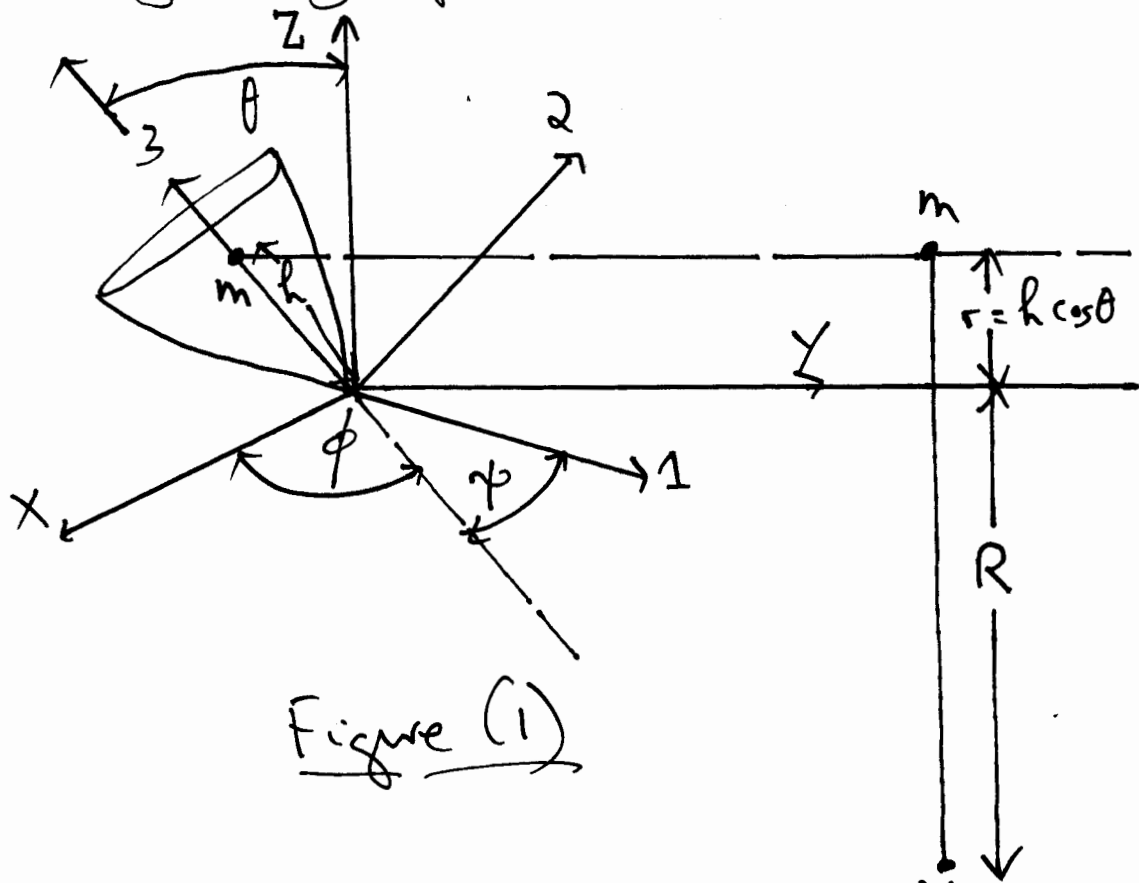


Figure (1)

In Fig (1) the point of the gyro is allowed to move in the gravitational field of the earth of mass M and radius R .

The Lagrangian is therefore made up of translational and rotational motions:

$$L = T_{\text{trans}} + T_{\text{rot}} - U \quad (1)$$

where T_{trans} is the translational kinetic energy and T_{rot} is the rotational kinetic energy. The potential energy is

$$U = mg(h \cos \theta + R) \quad (2)$$

where R is the radius of the earth and

$$r = l \cos \theta. \quad (3)$$

Here l is the constant distance from the origin to the centre of mass of the gyroscope along the z axis of frame $(1, 2, 3)$. This is the frame of the principal moments of inertia of the gyroscope:

$$I_{12} = I_1 = I_2 \text{ and } I_3. \quad (4)$$

The translational kinetic energy is therefore:

$$T_{\text{trans}} = \frac{1}{2} \frac{mM}{m+M} \underline{v} \cdot \underline{v} \quad (5)$$

where:

$$\underline{v} = \frac{d}{dt} (\underline{r} + \underline{R}) = \dot{\underline{r}} \quad (6)$$

because:

$$\frac{dR}{dt} = 0. \quad (7)$$

Therefore:

$$T_{\text{trans}} = \frac{1}{2} m \dot{r}^2 \quad (8)$$

because $M \gg m$. $\quad (9)$

The rotational kinetic energy is as defined in

UFT 368:

$$T_{\text{rot}} = \frac{1}{2} I_{12} (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{1}{2} I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 \quad (10)$$

where Euler angles θ , ϕ and ψ are defined in Fig. (1).

3) The Lagrangian (1) is therefore:

$$L = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} I_{12} (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{1}{2} I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 - mg(r+R) \quad - (11)$$

in which

$$r = h \cos \theta \quad - (12)$$

There are four Euler Lagrange equations:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) \quad - (13)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) \quad - (14)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) \quad - (15)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\psi}} \right) = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\psi}} \right) \quad - (16)$$

Eqs. (15) and (16) give the same results as in
uFT 348:

$$\dot{\phi} = \frac{L_{\phi} - L_{\psi} \cos \theta}{I_{12} \sin^2 \theta} \quad - (17)$$

$$\dot{\psi} = \frac{1}{I_3} (L_{\psi} - I_3 \dot{\phi} \cos \theta) \quad - (18)$$

Eq. (14) gives:

$$\ddot{\theta} = \frac{\sin \theta}{I_{12}} \left(\dot{\phi}^2 \cos \theta (I_{12} - I_3) - I_3 \dot{\phi} \dot{\psi} + mgh \right) \quad - (19)$$

because r in the potential energy term of eq. (11) is defined by eq. (12). Eq. (19) is also the same as in UFT 368.

However, there is another equation (13) which

gives:
$$-mg = m\ddot{r} \quad - (20)$$

$$\ddot{r} = -g = \frac{MG}{R^2} \quad - (21)$$

i.e

This is the correct result for the motion of r if it is assumed that R is a constant. This assumes that the point of the gyro does not move w.r.t respect to the centre of the earth. More generally R depends on time. Defining:

$$\underline{R}_1 = \underline{r} + \underline{R} \quad - (22)$$

The kinetic energy is:

$$T = \frac{1}{2} m \underline{R}_1 \cdot \underline{R}_1 = \frac{1}{2} m R_1^2 \quad - (23)$$

and the potential energy is:

$$U = mgR_1 \quad - (24)$$

and
$$\frac{\partial \mathcal{L}}{\partial R_1} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{R}_1} \right) \quad - (25)$$

Eq. (25) means that the point of the gyroscope can move up or down.

3) The Lagrangian in Eq. (25) is

$$L = \frac{1}{2} m \dot{R}_1^2 + \frac{1}{2} I_2 (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{1}{2} I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 \quad - (26) \quad - \frac{m M G}{R_1}$$

in which

$$\begin{aligned} R_1 &= r + R \\ &= l \cos \theta + R \end{aligned} \quad - (27)$$

Eqs. (25) and (26) give:

$$\ddot{R}_1 = \frac{M G}{R_1^2} \quad - (28)$$

i.e.

$$\frac{d^2}{dt^2} (l \cos \theta(t) + R(t)) = \frac{M G}{(l \cos \theta + R)^2} \quad - (29)$$

i.e.

$$\ddot{R} + l \frac{d^2}{dt^2} (\cos \theta(t)) = \frac{M G}{(l \cos \theta + R)^2} \quad - (30)$$

Finally we:

$$\frac{d^2 \cos \theta(t)}{dt^2} = - (\ddot{\theta} \sin \theta + \dot{\theta}^2 \cos \theta) \quad - (31)$$

so

$$\boxed{\ddot{R} - l (\ddot{\theta} \sin \theta + \dot{\theta}^2 \cos \theta) = \frac{M G}{(l \cos \theta + R)^2}} \quad - (32)$$

Therefore eqs. (17) to (19) and (32) must be solved simultaneously to obtain $R(t)$.