

368(7): Investigation of Gyroscope Lift.

In the absence of a lifting force the Lagrangian is:

$$L = T - U \quad - (1)$$

where the rotational kinetic energy is:

$$T = \frac{1}{2} I_{12} (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2) + \frac{1}{2} I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 \quad - (2)$$

and the gravitational potential energy is:

$$U = mgh \cos \theta \quad - (3)$$

The Euler Lagrange equations of motion are:

$$\ddot{\theta} = \frac{\sin \theta}{I_{12}} \left(\dot{\phi}^2 \cos \theta (I_{12} - I_3) - I_3 \dot{\phi} \dot{\psi} + mgh \right) \quad - (4)$$

$$\dot{\phi} = \frac{L_{\phi} - L_{\psi} \cos \theta}{I_{12} \sin^2 \theta} \quad - (5)$$

and

$$\dot{\psi} = \frac{1}{I_3} (L_{\psi} - I_3 \dot{\phi} \cos \theta) \quad - (6)$$

which have been solved numerically by Horst Eckardt to give rotation and precession in a well defined way. This is the first complete solution of gyroscope motion since "Mechanique Analytique" (1811 and 1815) by Lagrange.

The simplest way to investigate lift in a gyroscope is to decrease the effective

mass: $m \rightarrow m - m_1 - (7)$

so eq. (4) becomes:

$$\ddot{\theta} = \frac{\sin \theta}{I_{12}} \left(\dot{\phi}^2 \cos \theta (I_{12} - I_3) - I_3 \dot{\phi} \dot{\psi} + (m - m_1) g h \right) - (8)$$

Eqs. (5) and (6) remain the same.

Therefore solve the equations (8) and (5) and (6) simultaneously to find the effect of m_1 on the nutation and precession of the gyroscope.

The gyroscope appears to be weightless when:

$$m = m_1 - (9)$$

so under the weightless condition:

$$\ddot{\theta} = \frac{\sin \theta}{I_{12}} \left(\dot{\phi}^2 \cos \theta (I_{12} - I_3) - I_3 \dot{\phi} \dot{\psi} \right) - (10)$$

$$\dot{\phi} = \frac{L \dot{\phi} - L \dot{\psi} \cos \theta}{I_{12} \sin^2 \theta} - (11)$$

$$\dot{\psi} = \frac{1}{I_3} (L \dot{\psi} - I_3 \dot{\phi} \cos \theta) - (12)$$

Solving eqs. (10) to (12) gives a well defined nutation and precession, the nutation of a weightless gyroscope.