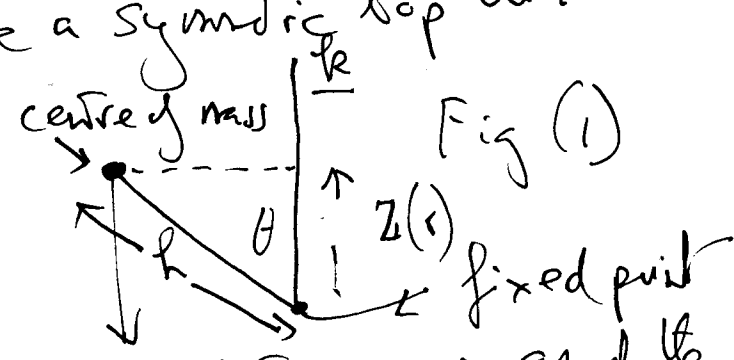


3/8(2) : The Lagrangian of a Gyroscope in Spherical Polar Coordinates

Consider the gyroscope to be a symmetric top with one part fixed, with mass  $m$ .  
 Will refer to Fig (1):



$$z = h \cos \theta \quad - (1)$$

so the potential energy of attraction between  $m$  and the mass of the earth  $M$  is

$$U = mgh \cos \theta \quad - (2)$$

In spherical polar coordinates the linear velocity is:

$$\underline{v} = \dot{r} \underline{e}_r + r \dot{\theta} \underline{e}_\theta + r \dot{\phi} \sin \theta \underline{e}_\phi \quad - (3)$$

and the angular momentum is:

$$\underline{L} = m \underline{r} \times \underline{v} = m r^2 (\dot{\theta} \underline{e}_\phi - \dot{\phi} \sin \theta \underline{e}_\theta) \quad - (4)$$

From eq. (4):  $L^2 = m^2 r^4 (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) \quad - (5)$

is a constant of motion. The  $z$  component of angular

momentum is:

$$L_z = m r^2 \dot{\phi} \sin^2 \theta \quad - (6)$$

The Lagrangian is therefore:

$$\mathcal{L} = \frac{1}{2} m v^2 - U \quad - (7)$$

$$2) = \frac{1}{2} m (\dot{r}^2 + r^2 (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta)) - mgl \cos \theta \quad (8)$$

As in 4FT 270, define:

$$\dot{\beta}^2 = \dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta \quad (9)$$

so 
$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\beta}^2) - mgl \cos \theta \quad (10)$$

$I_{12}$  terms of Euler angles (Moria and Thomson Sec 10.10), the same Lagrangian is:

$$L = \frac{1}{2} I_{12} \dot{\beta}^2 + \frac{1}{2} I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 - mgl \cos \theta \quad (11)$$

so 
$$I_{12} = mr^2 \quad (12)$$

and 
$$I_3 (\dot{\phi} \cos \theta + \dot{\psi})^2 = mr \dot{r}^2 \quad (13)$$

Also: 
$$\dot{\beta}^2 = \omega_1^2 + \omega_2^2 = \dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2 \quad (14)$$

and 
$$\omega_3^2 = (\dot{\phi} \cos \theta + \dot{\psi})^2 = \frac{mr \dot{r}^2}{I_3} \quad (15)$$

so the Lagrangian is:

$$L = \frac{1}{2} I_3 \omega_3^2 + \frac{1}{2} I_{12} (\omega_1^2 + \omega_2^2) - mgl \cos \theta \quad (16)$$

and the kinetic energy is:

$$T = \frac{1}{2} I_3 \omega_3^2 + \frac{1}{2} I_{12} (\omega_1^2 + \omega_2^2) \quad (17)$$

3) The two Euler Lagrange equations are:

$$\frac{\partial \mathcal{L}}{\partial r} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}} \quad - (18)$$

and

$$\frac{\partial \mathcal{L}}{\partial \beta} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\beta}} \quad - (19)$$

Eqs. (10) and (19) give:

$$m \ddot{r} = m r \dot{\beta}^2 - m g \frac{dz(r)}{dr} \quad - (20)$$

using

$$r = h \quad - (21)$$

in eq. (10).

Eqs. (10) and (19) give:

$$\frac{\partial \mathcal{L}}{\partial \beta} = 0, \quad \frac{\partial \mathcal{L}}{\partial \dot{\beta}} = m \dot{\beta} r^2 \quad - (22)$$

and

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\beta}} \right) = 0 \quad - (23)$$

The angular momentum:

$$L = \frac{\partial \mathcal{L}}{\partial \dot{\beta}} = m \dot{\beta} r^2 \quad - (24)$$

is a constant of motion:

$$\frac{dL}{dt} = 0 \quad - (25)$$

4) Here:

$$\dot{\beta} = (\omega_1^2 + \omega_2^2)^{1/2} = (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta)^{1/2} \quad - (26)$$
$$= \frac{L}{mr^2}$$

From eqs. (20) and (26):

$$m\ddot{r} = \frac{L^2}{mr^3} - mg \frac{dZ(r)}{dr} \quad - (27)$$

From eq. (15):

$$m\dot{r}^2 = I_3 \omega_3^2 \quad - (28)$$

From eq. (26):

$$\omega_1^2 + \omega_2^2 = \frac{L^2}{mr^4} \quad - (29)$$

The Hamiltonian:

$$H = \frac{1}{2} I_3 \omega_3^2 + \frac{1}{2} I_{10} (\omega_1^2 + \omega_2^2) + mgZ(r) \quad - (30)$$

is also a constant of motion:

$$\frac{dH}{dt} = 0 \quad - (31)$$

These are the equations of motion of the gyro developed as a symmetric top with one point fixed.

From eqs. (28), (29) and (30):

$$H = \frac{1}{2} m \dot{r}^2 + \frac{I_1 \omega^2}{2mr^4} + mgz(r) \quad (32)$$

and from eq. (27):

$$m \ddot{r} = \frac{L^2}{mr^3} - mg \frac{dz(r)}{dr} \quad (33)$$

with

$$\frac{dH}{dt} = 0 \quad (34)$$

These are equations for the position  $z(r)$  of the gyro, the height of the centre of mass above the fixed point of the gyro.

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